

# All possible canonical second-order three-impedance class-A and class-B oscillators

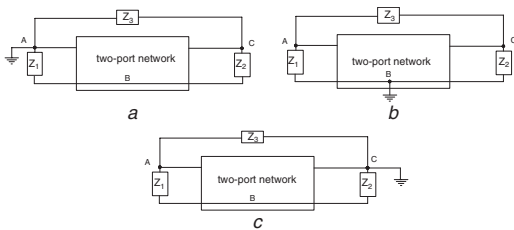
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A report is presented on all the possible canonical second-order single-transistor three-impedance oscillators that can be derived from the general class-A or class-B two-port network topologies. Two of the ten reported cases are already known as the class-A and class-B Colpitts oscillators. A selected case, which has the appealing feature of extending the oscillation frequency beyond the resonance frequency, is chosen for experimental validation.

**Introduction:** Two-port network modelling offers a high level of abstraction which enables device-independent circuit structure classification [1, 2]. In a recent publication [3], single-transistor three-impedance oscillators were classified as class-A, class-B or class-C using abstract two-port network modelling. Fig. 1 shows all possible three-impedance topologies that can be connected to a two-port network. Based on which terminal is grounded, the three topologies are classified as class-A, class-B or class-C, where the terminals A, B and C are marked within Fig. 1. It was particularly shown in [3] that class-C cannot yield an oscillator whether it is implemented using a BJT or an MOS transistor while the class-A and class-B topologies have two distinct characteristic equations given respectively by

$$Z_1 + Z_2 + Z_3 - g_m Z_1 Z_3 = 0 \text{ class-A} \quad (1a)$$

$$Z_1 + Z_2 + Z_3 + g_m Z_1 Z_2 = 0 \text{ class-B} \quad (1b)$$



**Fig. 1** Three possible topologies for three-impedance two-port network structure

a Class-A  
b Class-B  
c Class-C

where  $g_m$  is the small-signal transconductance of a BJT or an MOS transistor. More sophisticated versions of (1), which take into consideration parasitic capacitances and other non-ideal effects, where derived in [3] based on general two-port transmission parameters. However, equations (1a and b) are the starting point for oscillator circuit design by properly choosing the impedances  $Z_{1,2,3}$  such that there exists a positive value for  $g_m$  to start-up oscillations and result in a positive oscillation frequency  $\omega_o$ .

Here we report all the valid possibilities of  $Z_1, Z_2$  and  $Z_3$  which result in a canonical second-order oscillator for both class-A and class-B. A canonical second-order circuit employs a maximum of two resistors. The constraint in [3] that two capacitors and one inductor must be used, which is relevant only to the Colpitts oscillator structure, is lifted while the second-order constraint is imposed.

**Table 1:** All possible cases for class-A topology which yield a canonical second-order oscillator

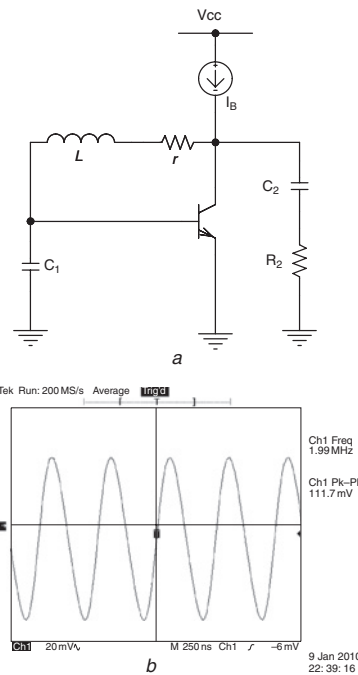
Case	$Z_1$	$Z_2$	$Z_3$	$g_m$	$\omega_o^2$
1	$1/sC_1$	$1/sC_2$	$r + sL$	$\frac{r}{L/C_1}$	$\frac{1}{LC_{eff}} - \frac{r^2}{L^2}$
2	$1/sC_1$	$1/sC_2$	$R \parallel (r + sL)$	$\frac{r}{L/C_1} + \frac{1 + (C_1/C_2)}{R}$	$\frac{1}{LC_{eff}} - \frac{r^2}{L^2}$
3	$1/sC_1$	$R + 1/sC_2$	$r + sL$	$\frac{r + R}{L/C_1}$	$\frac{1}{LC_{eff}} - \frac{r(r + R)}{L^2}$
4	$1/sC_1$	$R$	$r + sL$	$\frac{r + R}{L/C_1}$	$\frac{1}{LC_1} - \frac{r(r + R)}{L^2}$
5	$r + sL$	$1/sC_2$	$R + 1/sC_3$	$\frac{r + R}{(L/C_3) + rR}$	$\frac{1}{LC_{eff}} \left( \frac{L + rRC_3}{L - R^2 C_3} \right) - \frac{r(r + R)}{L^2 - LR^2 C_3}$
6	$R_1 \parallel (1/sC_1)$	$1/sC_2$	$R_3$	$\frac{1}{R_3} + \frac{1 + (C_1/C_2)}{R_1}$	$\frac{1}{R_1 R_3 C_1 C_2}$

**Canonical class-A oscillators:** A Matlab code was written to symbolically find all valid characteristic equations of this class and eliminate valid third-order cases as well as non-canonical second-order ones. The remaining are 12 valid circuits. Noting from (1a) that  $Z_1$  and  $Z_3$  are interchangeable means that effectively there are only six different circuits, as summarised in Table 1. The ideal start-up value of  $g_m$ , which corresponds to a Hopf-type bifurcation and the appearance of a pure imaginary eigen-pair, is given in Table 1 alongside the oscillation frequency  $\omega_o$ .

It can be seen that  $\omega_o$  for the first three cases is a function of  $C_{eff} = C_1 C_2 / (C_1 + C_2)$ . However, only the first and third cases have a start-up condition independent of  $C_2$  allowing for  $\omega_o$  to be independently tuned through  $C_2$  without affecting the start-up condition. It is also clear that cases 4 and 6 are the only two cases which employ two reactive elements and are truly second-order oscillators while the other cases employ three reactive elements and are degenerate second-order systems. Case 5 is clearly not a practical case owing to the complex expression of  $\omega_o$ . In conclusion, the practically appealing circuits seem to be those of cases 1 and 3 if independent frequency tuning is required and cases 4 and 6 if only two reactive elements are to be used. Note that case 1 is already known as the class-A Colpitts oscillator [3].

**Table 2:** All possible cases for class-B topology which yield a canonical second-order oscillator

Case	$Z_1$	$Z_2$	$Z_3$	$g_m$	$\omega_o^2$
1	$1/sC_1$	$1/sC_2$	$r + sL$	$\frac{r}{L/(C_1 + C_2)}$	$\frac{1}{LC_{eff}}$
2	$1/sC_1$	$1/sC_2$	$R \parallel (r + sL)$	$\frac{(r + R)(C_1 + C_2 + rRC_1 C_2/L)}{R^2 C_{eff} - L}$ ; $L < R^2 C_{eff}$	$\frac{(1 + r/R)^2}{LC_{eff} - (L/R)^2}$
3	$1/sC_1$	$R \parallel (1/sC_2)$	$r + sL$	$\frac{r + R}{R^2 C_2 / C_1} + \frac{r}{L/(C_2 + C_1 + rC_1/R)}$	$\frac{1}{LC_{eff}} + \frac{r/R}{LC_2}$
4	$1/sC_1$	$R + 1/sC_2$	$r + sL$	$\frac{C_1 + C_2}{L/(r + R)} - RC_2$ ; $L = mR(r + R)C_2$ ; $m > 1$	$\frac{m}{m - 1} \frac{1}{LC_{eff}}$



**Fig. 2** Circuit realising case 4 of Table 2, and experimental waveform across  $R_2$

a Circuit realising case 4 of Table 2  
b Experimental waveform across  $R_2$

**Canonical class-B oscillators:** Using a Matlab code to evaluate the different possibilities of this class and noting from (1b) that  $Z_1$  and  $Z_2$  can be symmetrically interchanged, only four canonical second-order circuits were found, as listed in Table 2. Case 1 is already known as the class-B Colpitts oscillator [3]. Note from Table 2 that case 2 and case 4 impose a constraint on the value of L. However, case 3 and case 4 have an appealing feature where  $\omega_o$  can be made higher than the resonance frequency  $1/\sqrt{LC_{eff}}$ . This is not possible for any of the

class-A cases in Table 1. In conclusion, the two most practical cases of this class are cases 1 and 4.

*Experimental validation:* The circuit shown in Fig. 2a was constructed to validate case 4 in Table 2 as an example. The components chosen were:  $L = 116 \mu\text{H}$ ;  $r = 140 \Omega$  ( $107 \Omega$  physical resistance +  $33 \Omega$  internal inductor resistance),  $R_2 = 1 \text{ k}\Omega$ ;  $C_1 = C_2 = 100 \text{ pF}$  and a Q2N2222 NPN transistor was used. These values yield  $m = 2.035$  and a resonance frequency  $f_{\text{res}} = 1/2\pi\sqrt{LC_{\text{eff}}} = 1.477 \text{ MHz}$ . Consulting Table 2, it is expected that the measured oscillation frequency should be equal to  $\sqrt{m/(m-1)}f_{\text{res}} \simeq 2.07 \text{ MHz}$ . Fig. 2b shows the observed experimental waveform across  $R_2$  with a measured oscillation frequency of  $1.99 \text{ MHz}$  very close to the expected value. The bias current to start-up oscillations was  $I_B = 53 \mu\text{A}$ . Note that taking the output across  $R_2$  has the advantage of being free from the DC biasing voltage of the employed transistor.

*Conclusion:* Out of 95 possible class-A oscillators and 19 possible class-B oscillators, only six class-A and four class-B were found to yield a canonical second-order oscillator and are reported here. Most of the found circuits are second-order degenerate structures which, with non-linear modelling, will remain third-order circuits [4].

However, the characteristic equations (1a and b) are valid only for small-signal linearised transistor models. As such, the second-order classification remains valid.

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14 January 2010  
doi: 10.1049/el.2010.0082

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